

Equilibria configurations for epitaxial crystal growth with adatoms

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Joint work with Marco Caroccia and Laurent Dietrich



BIRS Workshop

Topics in the Calculus of Variations: Recent Advances and New Trends

May 21, 2018

Outline of the talk

- ▶ Surface evolution

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- ▶ A model with adatoms

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- ▶ The big plan

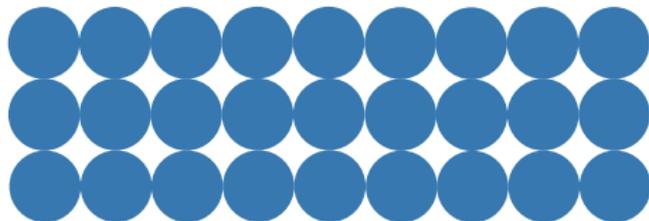
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- ▶ Surface evolution
- ▶ A model with adatoms
- ▶ The big plan
- ▶ What we have done/are doing

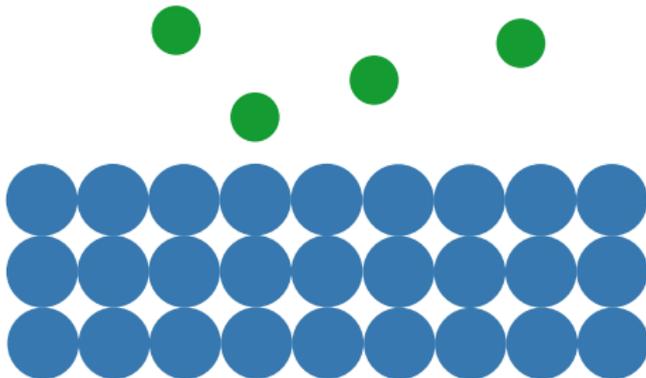
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- ▶ Surface evolution
- ▶ A model with adatoms
- ▶ The big plan
- ▶ What we have done/are doing
- ▶ Future plans

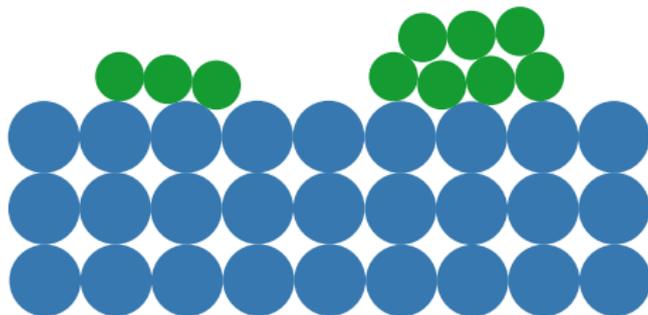
Epitaxial growth



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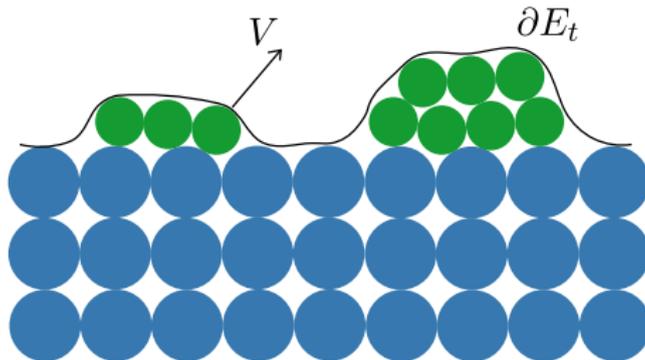


Epitaxial growth - Surface evolution

According to the *Einstein-Nernst law*, the surfaces $\{E_t\}_{t \geq 0}$ evolve following the volume preserving equation

$$\rho V = D \Delta_{\partial E_t} \mu \quad \text{on } \partial E_t,$$

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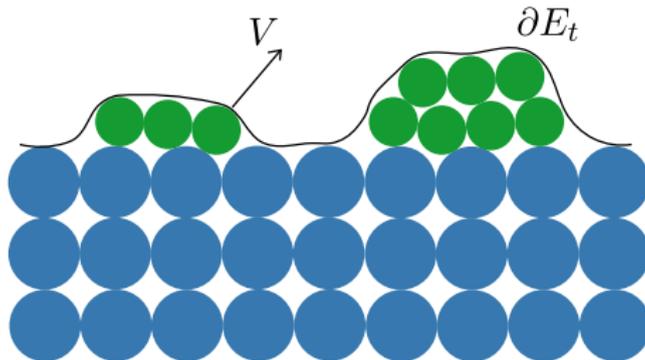
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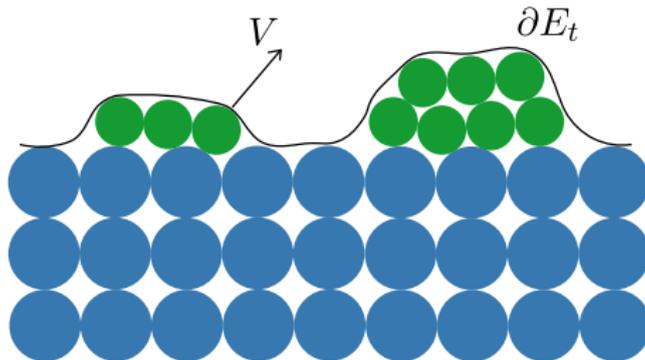
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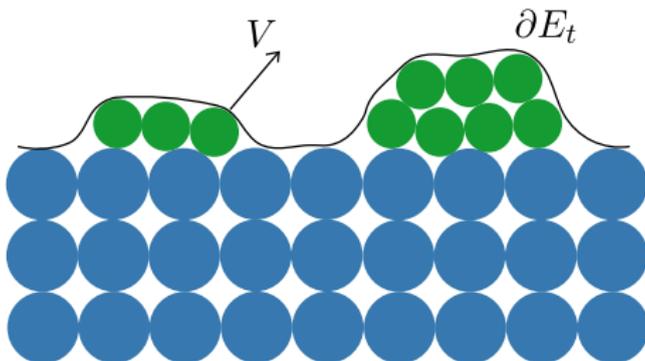
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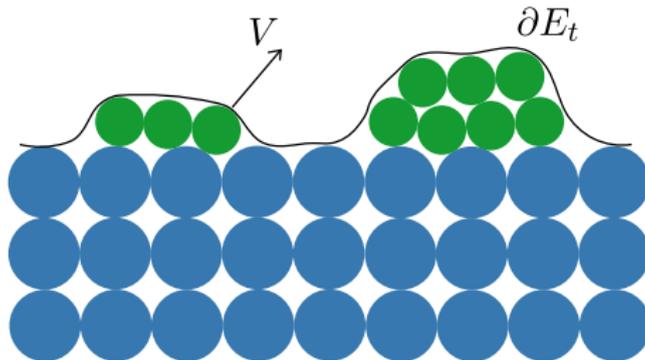
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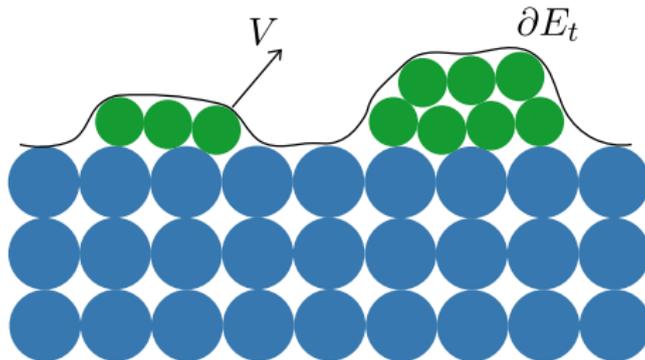
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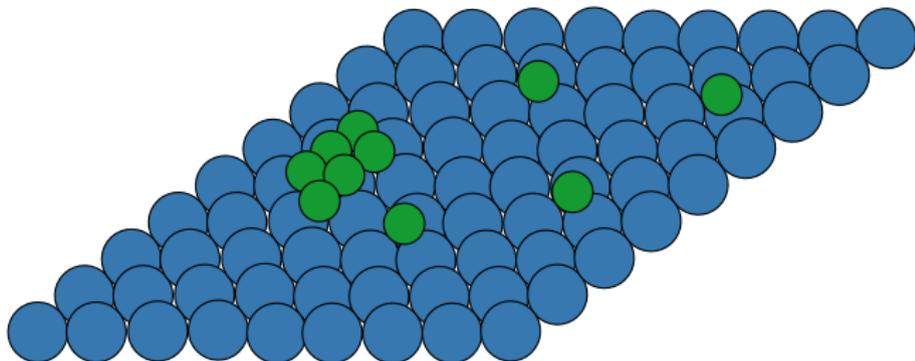


Adatoms

On the surface:

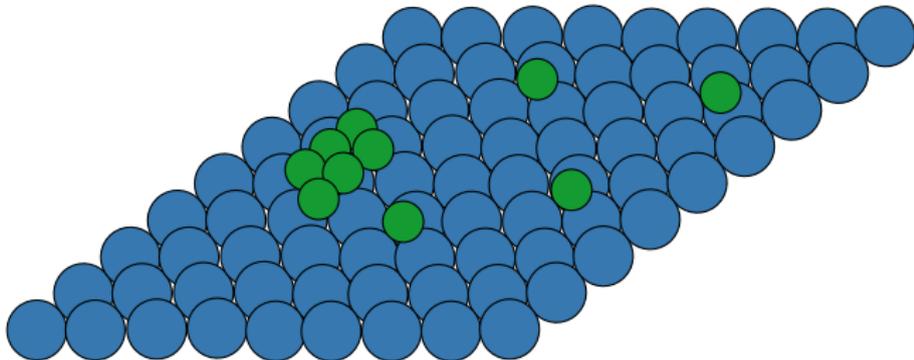
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On the surface: Atoms



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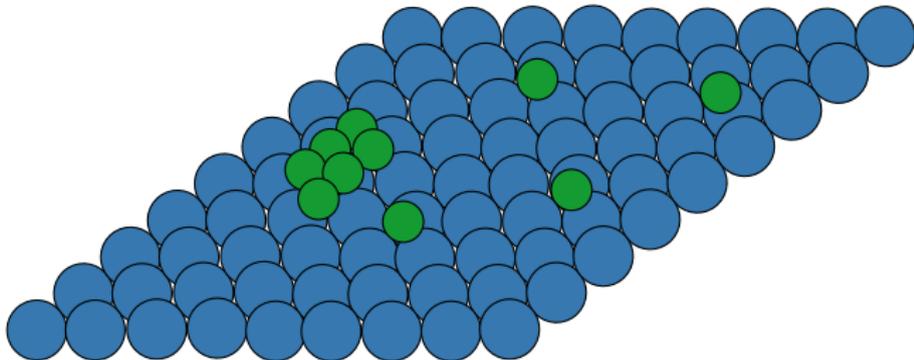
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Why consider adatoms?

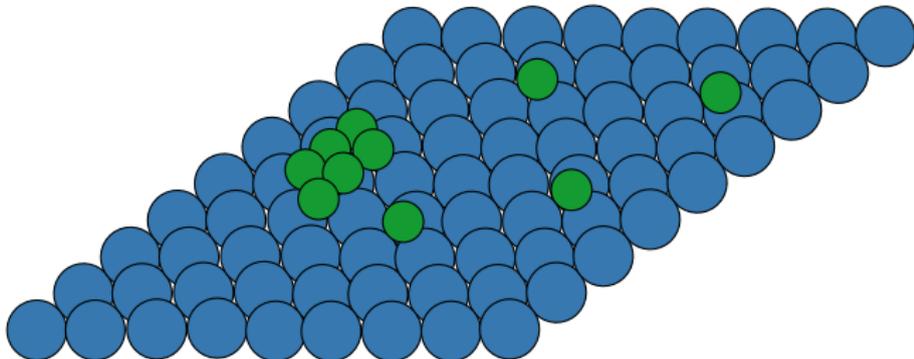


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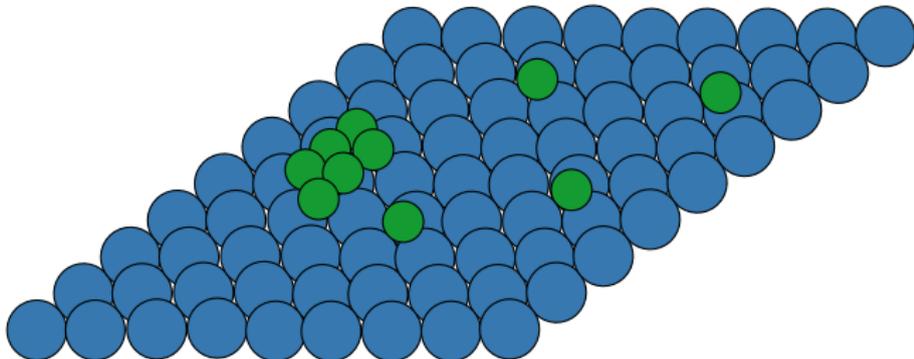
- ▶ important in models for solid-vapor interfaces



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Why consider adatoms?

- ▶ important in models for solid-vapor interfaces
- ▶ effect of regularizing the unstable parabolic equations for surface evolution



A model with adatoms

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Eliot Fried and Morton E. Gurtin,

A unified treatment of evolving interfaces accounting for small deformations and atomic transport with emphasis on grain-boundaries and epitaxy,

Advances in applied mechanics, 40 (2004), pp. 1-177

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- ▶ the kinetic term bV : originates from the constitutive equation $F = bV$, where F is a dissipative force associated with the attachment of vapor atoms on the solid surface.

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- ▶ obtain the above system of equations as a gradient flow of some energy \mathcal{F} ✓
- ▶ study the energy \mathcal{F} ✓
- ▶ consider an approximation of \mathcal{F} via a diffuse interface energy \mathcal{F}_ε LOADING, PLEASE WAIT 
- ▶ study the two systems of evolution equations
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Starting from the standard evolutions equations $\rho V = D\Delta_{\partial E_t}\mu$ on ∂E_t , Fried and Gurtin model for surface diffusion including adatoms

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Previous works

Martin Burger,
Surface diffusion including adatoms,
Commun. Math. Sci., 4 (2006), pp. 1-51

Andreas Rätz, Axel Voigt,
A diffuse-interface approximation for surface diffusion including adatoms
Nonlinearity, 20 (2007), pp. 177-192

Christina Stöcker, Axel Voigt,
A level set approach to anisotropic surface evolution with free adatoms,
SIAM Journal on Applied Mathematics, 69 (2008), pp. 64-80

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We define the energy functional

$$\mathcal{F}(E, u) := \int_{\partial^* E} \psi(u) \, d\mathcal{H}^{N-1},$$

where $E \subset \mathbb{R}^N$ is a set of finite perimeter and $u \in L^1(\partial^* E; \mathbb{R}_+)$.

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- ▶ E is the *solid*
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- ▶ prototype of ψ is $\psi(s) := 1 + s^2/2$ (suggested by Fried and Gurtin)

The minimum problem

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Given $E \subset \mathbb{R}^N$ is a set of finite perimeter and $u \in L^1(\partial^* E; \mathbb{R}_+)$, we define the *total mass*

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Definition

For $m > 0$, we define the *admissible class of competitors*

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We are interested in the following *constrained* minimization problem

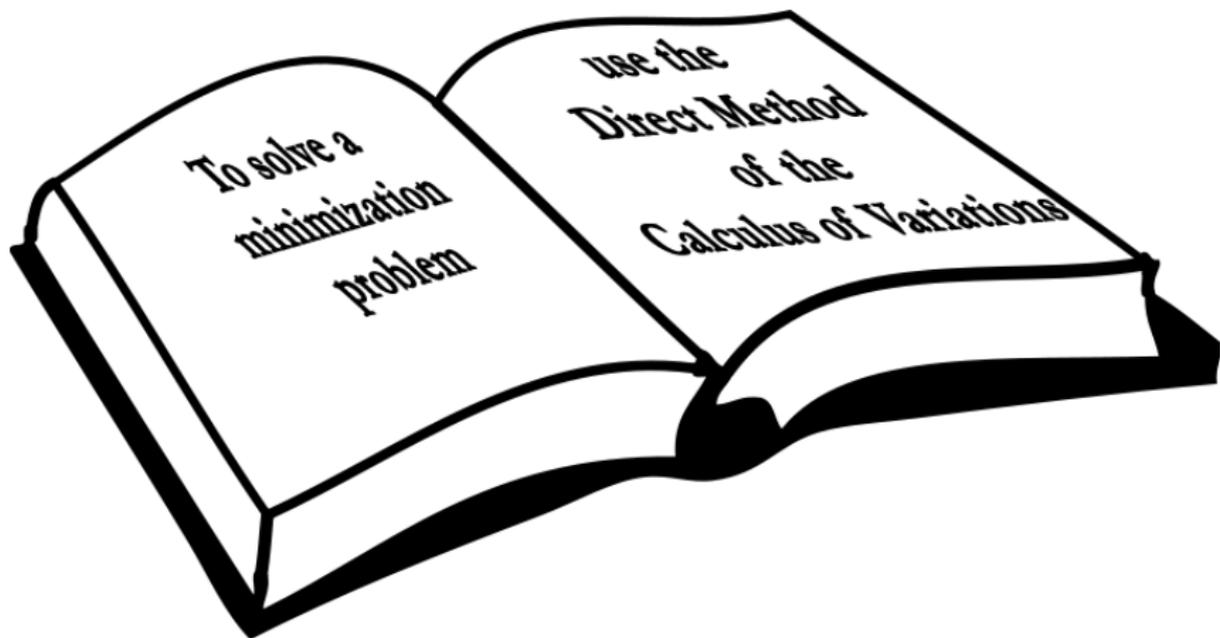
$$\gamma_m := \inf \left\{ \mathcal{F}(E, u) = \int_{\partial^* E} \psi(u) \, d\mathcal{H}^{N-1} : (E, u) \in \text{Cl}(m) \right\},$$

Equilibria configurations

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$$\begin{aligned}\mathcal{F}(E, u) &= \int_{\partial^* E} \psi(u) \, d\mathcal{H}^{N-1} \\ &\geq \mathcal{H}^{N-1}(\partial^* E) \psi \left(\frac{1}{\mathcal{H}^{N-1}(\partial^* E)} \int_{\partial^* E} u \, d\mathcal{H}^{N-1} \right)\end{aligned}$$

by applying Jensen's inequality

Equilibria configurations

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Equilibria configurations

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Achtung! Volume constraint

$$\mathcal{M}(E, u) = \rho|E| + \bar{u}\mathcal{H}^{N-1}(\partial^* E) = m.$$

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Fix $m > 0$. Assume ψ behaves nicely at $s = 0$ and at infinity (technical conditions).

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Remark

ψ behaves nicely at $s = 0$ and at infinity is in order to avoid as minimizers balls with zero radius and infinite adatom density or balls with zero adatom density.

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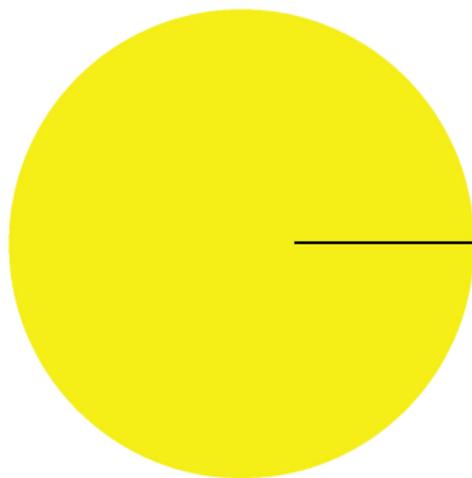
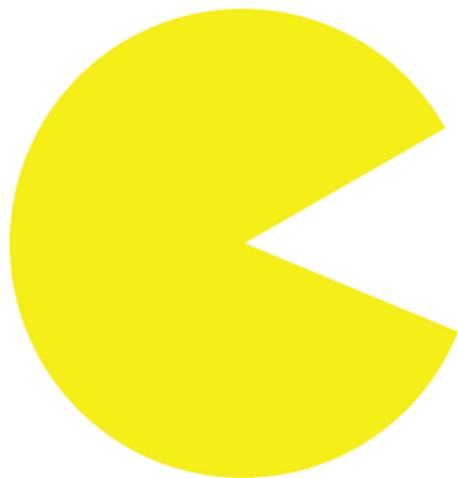
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Remark

Non-uniqueness of the solution: different size of balls (other than translation invariance).

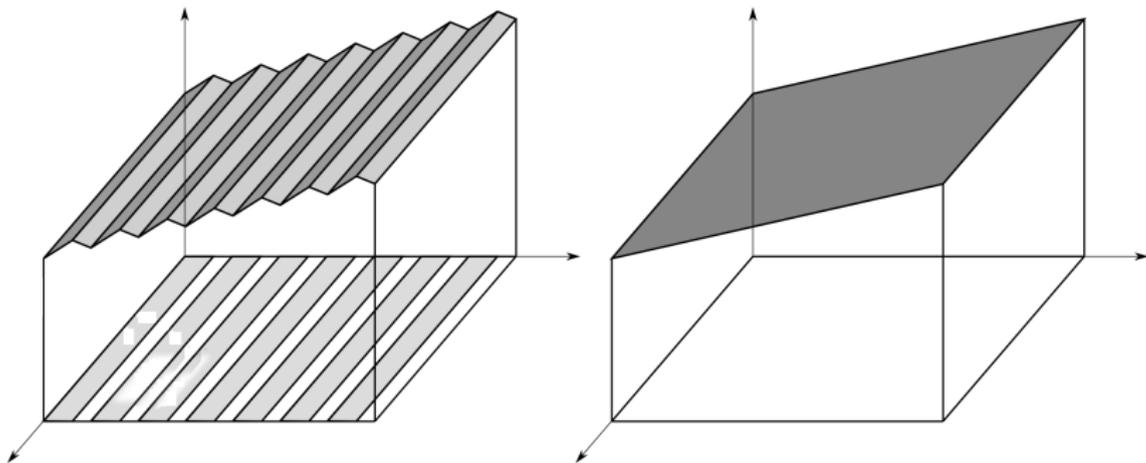
Is the energy l.s.c.?

Is the energy I.s.c.?



The pacman example

Is the energy l.s.c.?



The wiggling example

The extended energy

$$(E, u)$$

E set of finite perimeter

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$$\mathcal{F}(E, \mu) := \begin{cases} \int_{\partial^* E} \psi(u) \, d\mathcal{H}^{N-1} & \text{if } \mu = u |D\mathbb{1}_E| \text{ with } u \in L^1(\partial^* E; \mathbb{R}_+), \\ +\infty & \text{otherwise,} \end{cases}$$

where E is a set with finite perimeter

μ is a non-negative finite Radon measure on \mathbb{R}^N

in brief $(E, \mu) \in \mathfrak{G}$.

The topology

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We say that $((E_k, \mu_k))_{k \in \mathbb{N}} \subset \mathfrak{S}$ *convergence* to $(E, \mu) \in \mathfrak{S}$ if

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- ▶ $\mu_k \xrightarrow{*} \mu$ locally weakly*, i.e., for every $\varphi \in C_c(\mathbb{R}^N)$ we have that

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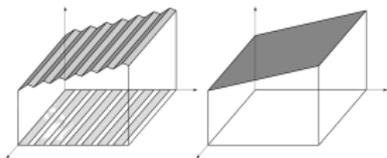
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Lemma

The above topology is metrizable.

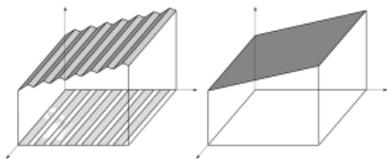
Necessary conditions for l.s.c.

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oscillating phenomena

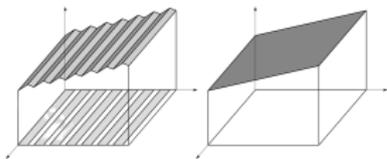
Necessary conditions for l.s.c.



ψ convex and subadditive

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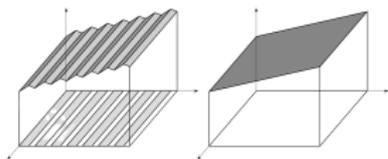
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concentration phenomena

Necessary conditions for I.s.c.



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recession function

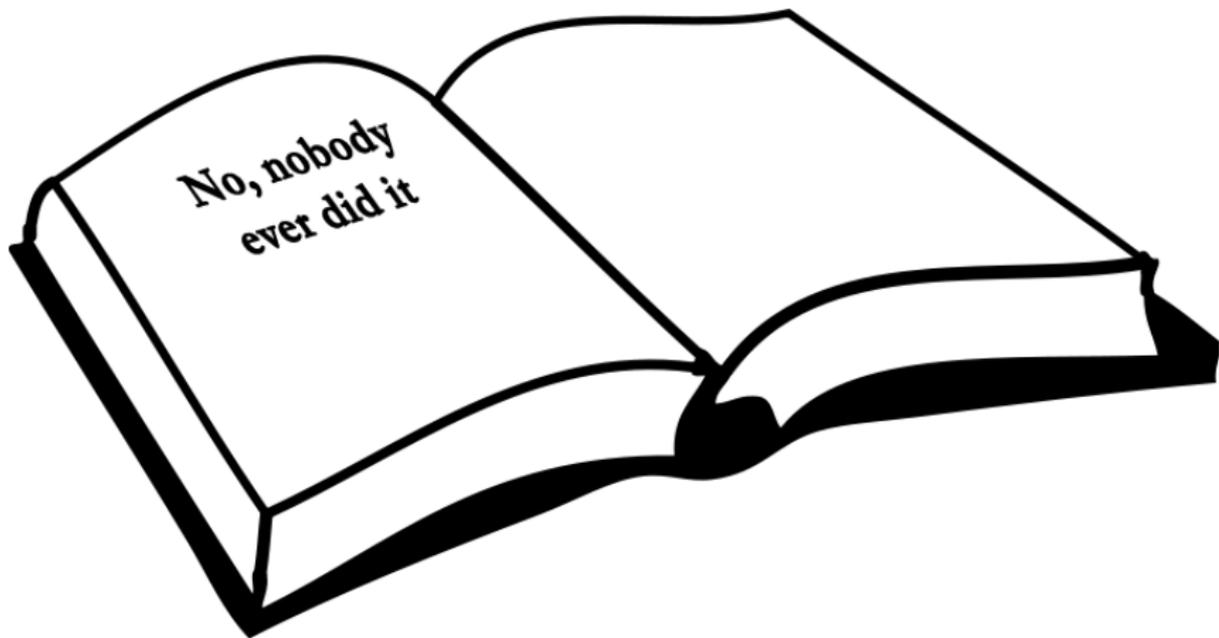
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Did anybody computed the relaxation of \mathcal{F} ?

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Convex subadditive envelope

Definition

Let $\psi : \mathbb{R} \rightarrow \mathbb{R}$. We say that ψ is *subadditive* if for every $r, s \in \mathbb{R}$,

$$\psi(r + s) \leq \psi(r) + \psi(s).$$

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Definition

Let $\psi : [0, \infty) \rightarrow \mathbb{R}$ be a function. We define its *convex subadditive envelope* $\bar{\psi} : [0, \infty) \rightarrow \mathbb{R}$ as

$$\bar{\psi}(s) := \sup\{ f(s) : f : [0, \infty) \rightarrow \mathbb{R} \text{ is convex, subadditive and } f \leq \psi \}.$$

Convex subadditive envelope

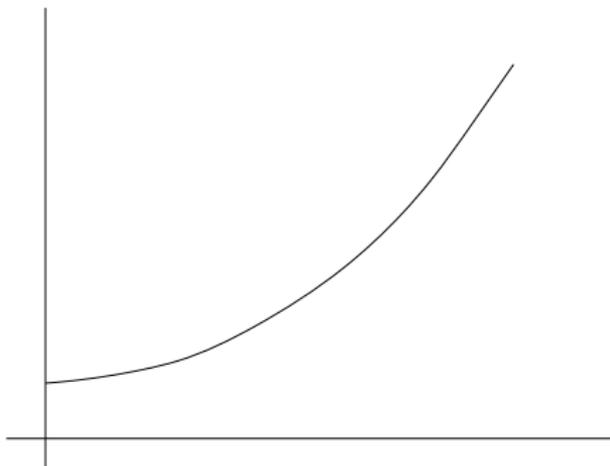
Lemma

$$\bar{\psi}(s) = \sup\{a_j s + b_j : j \in \mathbb{N}, b_j \geq 0\}.$$

Convex subadditive envelope

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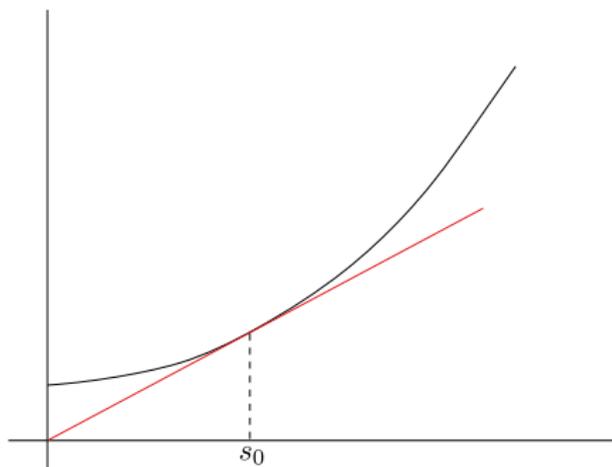
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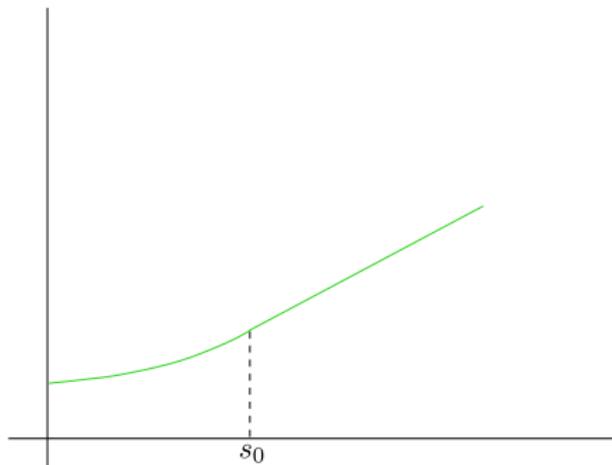
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Let $\psi : \mathbb{R}_+ \rightarrow (0, +\infty)$ be a C^1 convex non-decreasing function, and set

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We define the functional $\bar{\mathcal{F}} : \mathfrak{S} \rightarrow [0, \infty)$ as

$$\bar{\mathcal{F}}(E, \mu) := \int_{\partial^* E} \bar{\psi}(u) \, d\mathcal{H}^{N-1} + \Theta \mu^s(\mathbb{R}^N),$$

where we write $\mu = u\mathcal{H}^{N-1} \llcorner \partial^* E + \mu^s$ using the Radon-Nikodym decomposition.

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Theorem

$\bar{\mathcal{F}}$ is the relaxed functional of \mathcal{F} w.r.t. the topology in \mathfrak{G} .

Relaxation - liminf inequality

Liminf inequality:

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where $\mu_k := u_k \mathcal{H}^{N-1} \llcorner \partial^* E$.

By the l.s.c. of $\bar{\mathcal{F}}$ we obtain that

$$\liminf_{k \rightarrow \infty} \mathcal{F}(E_k, \mu_k) \geq \bar{\mathcal{F}}(E, \mu).$$

Relaxation - idea for the recovery sequence

Recovery sequence: let $(E, \mu) \in \mathfrak{G}$. Write

$$\mu = u\mathcal{H}^{N-1} \llcorner \partial^* E + \mu^s = u|D\mathbb{1}_E| + \mu^s,$$

using the Radon-Nikodym decomposition. Then

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We will construct:

(i) (F_k, v_k) with $(F_k, v_k) \rightarrow (E, u|D\mathbb{1}_E|)$ such that

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Relaxation - idea for the recovery sequence

Recovery sequence: let $(E, \mu) \in \mathfrak{G}$. Write

$$\mu = u\mathcal{H}^{N-1} \llcorner \partial^* E + \mu^s = u|D\mathbb{1}_E| + \mu^s,$$

using the Radon-Nikodym decomposition. Then

$$\overline{\mathcal{F}}(E, \mu) = \int_{\partial^* E} \overline{\psi}(u) \, d\mathcal{H}^{N-1} + \Theta\mu^s(\mathbb{R}^N).$$

We will construct:

(i) (F_k, v_k) with $(F_k, v_k) \rightarrow (E, u|D\mathbb{1}_E|)$ such that

$$\mathcal{F}(F_k, v_k) \rightarrow \int_{\partial^* E} \overline{\psi}(u) \, d\mathcal{H}^{N-1},$$

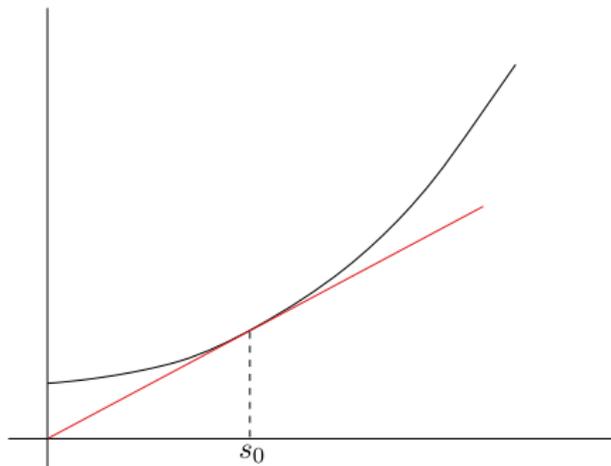
(ii) (G_k, w_k) with $(G_k, w_k) \rightarrow (\emptyset, \mu^s)$ such that

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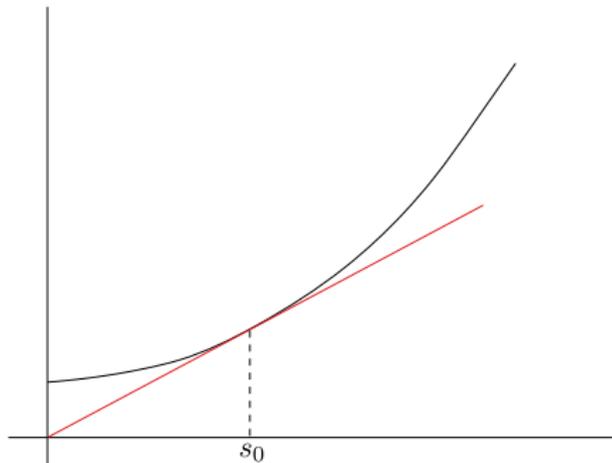
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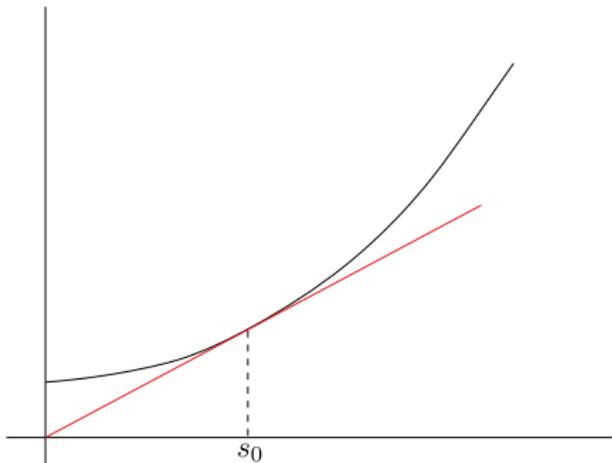
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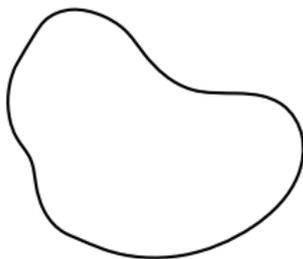
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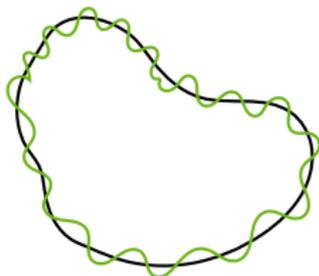
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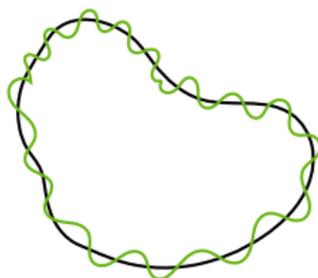
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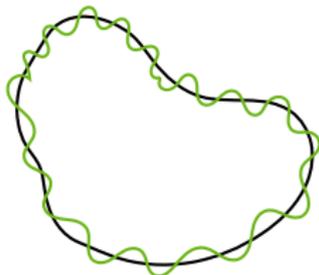
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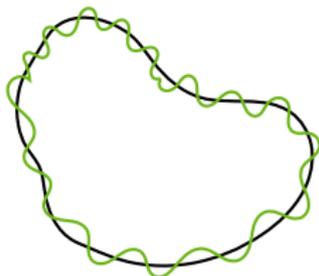
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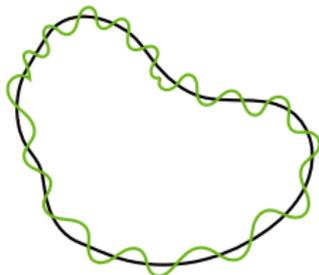
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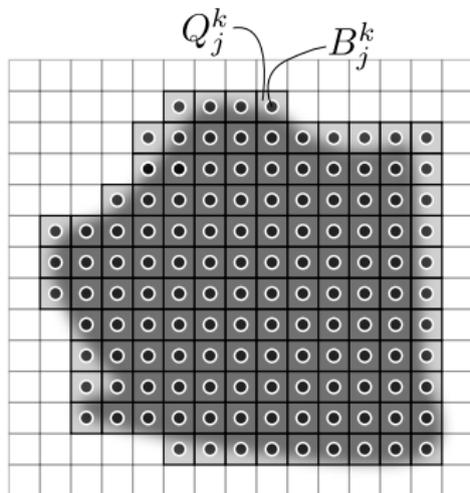
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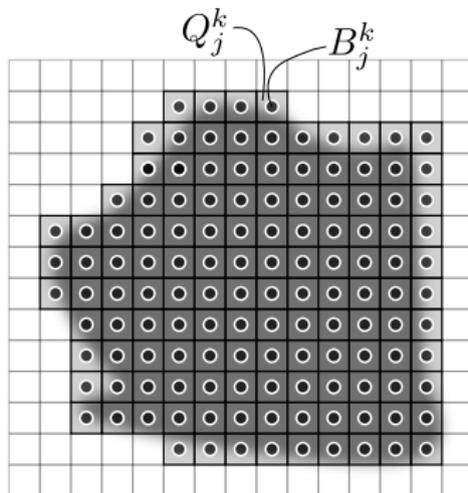
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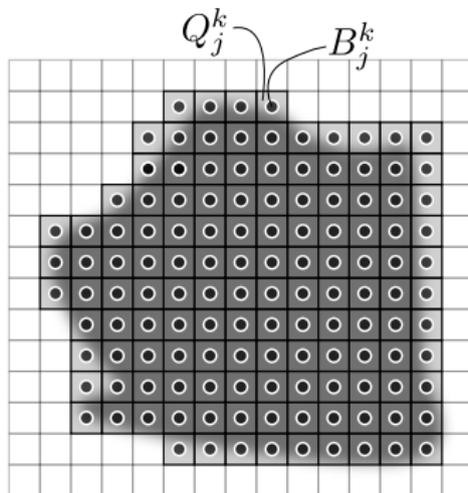
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Definition

Give $(E, \mu) \in \mathfrak{G}$, we define the *total mass*

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We are interested in the following constrained minimization problem

$$\overline{\gamma}_m := \inf \{ \overline{\mathcal{F}}(E, \mu) : (E, \mu) \in \overline{\text{Cl}}(m) \}.$$

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Remark

Due to the presence of the singular part of the measure, minimizers of $\overline{\mathcal{F}}$ have less structure than minimizers of \mathcal{F} .

A diffuse phase approximation of the energy

For $\varepsilon > 0$ define the *diffuse energy* $\mathcal{F}_\varepsilon : W^{1,2}(\mathbb{R}^N) \times C(\mathbb{R}^N) \rightarrow [0, +\infty]$ as

$$\mathcal{F}_\varepsilon(\phi, u) := \int_{\mathbb{R}^N} \left(\frac{1}{\varepsilon} W(\phi) + \varepsilon |\nabla \phi|^2 \right) \psi(u) dx.$$

where $W : \mathbb{R} \rightarrow \mathbb{R}$ is a double well potential.

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So that $\mathcal{F}_\varepsilon(\phi, u) \sim \mathcal{F}(E, u)$. Use the idea for the recovery sequence for $\overline{\mathcal{F}}$.

A discrete non-local approximation of the energy

Fix $R > 0$. For $n \in \mathbb{N}$ let $X_n := \{x_1, \dots, x_n\} \subset \mathbb{R}^N$ be such that x_i are chosen randomly in $B_R(0)$ uniformly.

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Theorem (True at 93% (2% more w.r.t. last week!))

If $\epsilon_n \rightarrow 0$ with a certain rate, then $\mathcal{F}_n^{(p)} \xrightarrow{\Gamma} c_{\eta,p,W} \overline{\mathcal{F}}$.

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- ▶ include more effects in the energy (more general materials)

That's all folks!

Thank you for your attention!

